

Efficient $\text{LaF}_3:\text{Nd}^{3+}$ -based vacuum-ultraviolet laser at 172 nm

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Vacuum-ultraviolet (VUV) laser radiation at 172 nm has been obtained from a solid-state $\text{LaF}_3:\text{Nd}^{3+}$ -based laser pumped by a pulsed-discharge molecular F_2 laser at 157 nm. The maximum slope efficiency of the solid-state laser described in this experiment was 21% (14% conversion efficiency), and the maximum output energy at 172 nm was 0.4 mJ for a nonoptimized optical cavity. This finding introduces serious prospects for realizing versions of active-medium-plus-source tunable VUV laser devices.

INTRODUCTION

Vacuum-ultraviolet (VUV) laser sources are useful devices for a variety of applications ranging from photochemistry to photolithography. However, until now only a few laser sources (such as H_2 and F_2), which have limited tunabilities, operate in this spectral region. Tunable VUV radiation has been obtained by a variety of methods, such as discharge excitation of atomic beams¹ and frequency mixing in gases and metal vapors,^{2,3} but tunable VUV radiation can in principle be obtained from a wide class of solid-state active media, such as rare-earth-activated wide-band-gap fluoride dielectric crystals because of the allowed radiative d - f transitions.⁴ Recent developments in new single-crystal wide-band-gap laser materials⁵⁻⁷ offer encouraging prospects for the development of crystal-based efficient VUV laser devices. Among the well-known fluoride crystals $\text{LaF}_3:\text{Nd}^{3+}$ ($\text{LaF}:\text{Nd}$) now seems to be the most suitable for VUV applications. The simplified energy-level diagram of Nd^{3+} ions in LaF_3 is shown in Fig. 1, and its excitation and emission spectra were studied previously.^{4,8-10} The strong and broad overlapping peaks of the excitation spectra (one of them centered at 159 nm) originate from transitions from the $4f^3$ ground state to the crystal field split $4f^25d$ levels of the Nd^{3+} ion in LaF_3 . However, among the five possible transitions that have been predicted by theoretical considerations concerning the C_2 site symmetry of the Nd^{3+} in LaF_3 , only four were observed.⁴ The observed VUV fluorescence originates from the $4f^25d \rightarrow 4f^3$ electric dipole-allowed transitions. They have been assigned to the $5d \rightarrow {}^4I_{11/2}$ and $5d \rightarrow {}^4I_{13/2}$ interconfigurational transitions.⁸ The fluorescence spectrum seems to have a simple structure at room temperature,⁹ and the VUV quantum yield of the Nd^{3+} ions in the LaF_3 host is more than 50%.⁸ The depression of the $4f^25d$ levels in LaF_3 caused by the crystal field is the lowest among the known fluorides, and the onset of the $4f^25d$ -configuration levels is observed at frequencies as high as $60\,000\text{ cm}^{-1}$. Therefore, the $\text{LaF}:\text{Nd}$

laser should have one of the shortest wavelengths of all possible Nd^{3+} -activated VUV lasers, provided that the appropriate pumping is available. Direct evidence of VUV laser action in $\text{LaF}:\text{Nd}$ was obtained recently,^{9,10} and this was the first solid-state laser in the VUV region of the spectrum. However, the method of excitation of the active medium that was used in Refs. 9 and 10 (the excitation of the active medium was due to fluorescence of Kr_2^* caused by an electron beam) was too complicated to be widely used at the laboratory scale. Therefore a simple and reliable pumping source could enhance the use of these sources for spectroscopic and other applications. Recently¹¹ we obtained laser action at 172 nm from $\text{LaF}_3:\text{Nd}^{3+}$ crystals pumped by a molecular F_2 laser at 157 nm with a 3% conversion efficiency. In this paper we report the observation of a 14% conversion efficiency with a 21% slope efficiency [defined as the $\tan(\Delta E_o/\Delta E_i)$, where ΔE_o is the change of the output energy when the input energy changes by ΔE_i near threshold] laser action at 172 nm from a $\text{LaF}:\text{Nd}$ crystal transversely pumped by a molecular F_2 laser at 157 nm. The energy conversion of this system is determined, and the processes that limit the energy available from the crystal are discussed.

EXPERIMENT

The $\text{LaF}:\text{Nd}$ crystals used in this experiment were grown from carbon crucibles by the Bridgman-Stockbarger method. Because the VUV transparency of the samples depends strongly on the amount and the types of impurities, during the growth process all possible methods of initial products and growth atmosphere refining were used. The laser samples, cylindrical rods 35 mm long with 5-mm radii, were pumped transversely, and the ends of the laser rods were optically flat and parallel with 10 s of arc. A polished flat window was made on the side of the cylindrical surface of the crystal for efficient excitation by the pumping laser source. The concentrations of the Nd^{3+} ions used in these crystals were between 0.3%

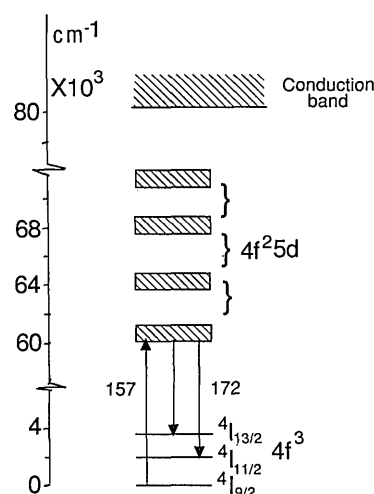


Fig. 1. Energy diagram of $\text{LaF}_3:\text{Nd}^{3+}$; the absorption at 157 nm and the laser emission transitions at 172 nm are indicated between the $4f^3$ and $4f^2 5d$ levels.

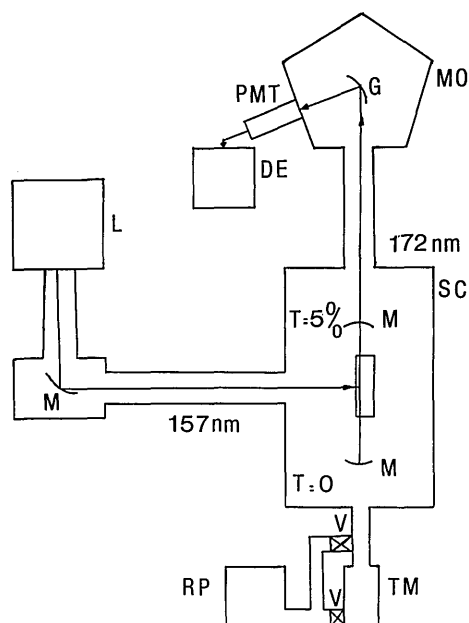


Fig. 2. Experimental layout. L, pulsed-discharge F_2 molecular laser at 157 nm; M's, mirrors; SC, stainless-steel vacuum chamber; MO, monochromator; G, grating; PMT, photomultiplier; DE, detection electronics; TM, turbo molecular pump; RP, rotary pump; V's, valves.

and 1%. The experimental apparatus is shown in Fig. 2. The pumping source was an F_2 molecular laser of the fast-discharge type that has been described previously.^{12,13} The maximum available energy from this device was 12 mJ at 3.5 atm (2.7 kTorr) total gas pressure. The pumping laser pulse had a width of 12 ns at half-maximum (Fig. 3). The pumping laser beam was focused on the crystal with an off-axis concave mirror with a 120-cm radius of curvature. The maximum energy density at the plane of the input window of the crystal was 120 mJ/cm² in a rectangular geometry of the focal spot with dimensions 0.3 mm \times 35 mm. The solid-state laser cavity was of the stable near-confocal type, and it was 15 cm long. Both the output and the back mirrors (Al + MgF_2 deposited on a substrate of BaF_2) of the cavity had a 30-cm

radius of curvature. The back mirror was totally non-transparent at 172 nm, while the transmittance of the output coupler at the same wavelength was 65%. The laser signal was detected with 0.2-m VUV monochromator (Acton VM 502), a solar-blind fast photomultiplier (EMI 9412), and a Tektronix 7104 oscilloscope. Because of strong absorption at these wavelengths by air, the optical path of the pumping and the solid-state laser beams was within stainless-steel vacuum lines pumped down to 10^{-5} mbar (7×10^{-6} Torr) with a turbomolecular pump. The pumping and the solid-state laser output energies were measured with a Gentek (PR100-type) pyroelectric detector.

RESULTS AND DISCUSSION

Stimulated emission from two $\text{LaF}_3:\text{Nd}$ crystals, pumped by an F_2 molecular laser at 157 nm, was obtained. The concentration of the Nd^{3+} dopant was 0.5 and 1.0 at. %. The pumping thresholds for obtaining stimulated emission from the two crystals were 3 and 0.5 mJ, respectively.

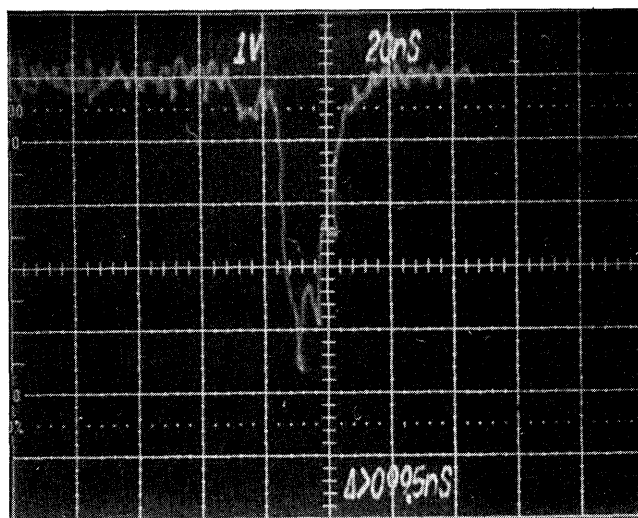


Fig. 3. Temporal evolution of the pumping F_2 laser pulse at 157 nm.

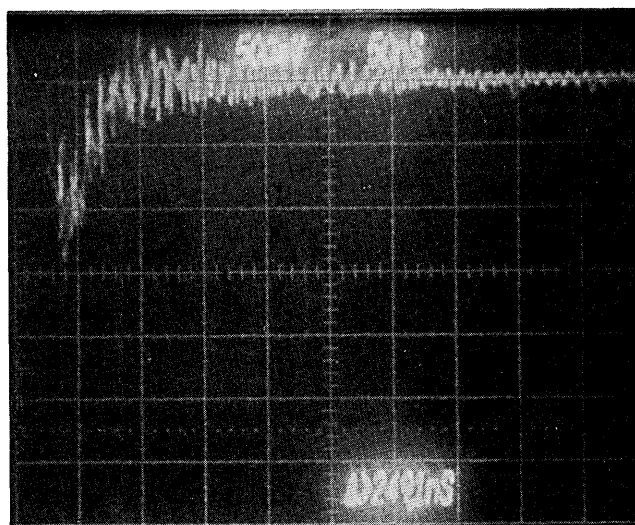


Fig. 4. Temporal evolution of the laser pulse from the $\text{LaF}_3:\text{Nd}^{3+}$ crystal at 172 nm.

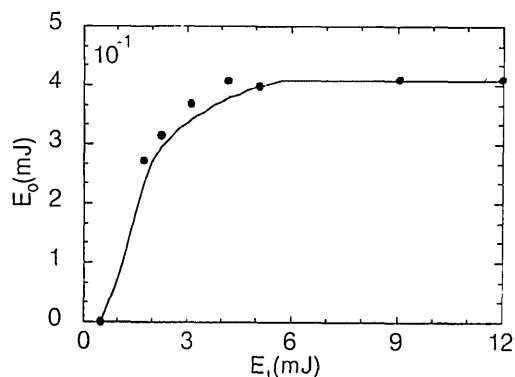


Fig. 5. Output energy E_o at 172 nm from the $\text{LaF}_3:\text{Nd}^{3+}$ crystal as a function of the input energy E_i of the pumping laser pulse at 157 nm.

Stimulated emission from a $\text{LaF}:\text{Nd}$ crystal was observed previously by Waynant^{9,10} for an only 0.1% concentration of the Nd^{3+} . The differences between our and Waynant's data are attributed to the different pumping geometries. The transverse pumping geometry of our experiment requires the active medium to have higher coefficients of absorption of the pumping radiation, which are unattainable at low doping levels. The temporal evolution of the generated pulse at 172 nm is indicated in Fig. 4. The pumping energy was 1.3 times threshold. Because the two pulses have similar temporal evolutions, the exhibition of high gain by the crystal is expected. The output energy of the $\text{LaF}:\text{Nd}$ VUV laser as a function of the pumping energy is indicated in Fig. 5. The slope efficiency of the described experimental setup for the linear part of the curve (Fig. 5) was 21%. Saturation is also taking place, while the mechanisms of this process are probably similar to those that occur for a $\text{YLF}:\text{Ce}^{3+}$ -laser,¹⁴ such as the formation of transient color centers by excited-state absorption of the pumping radiation.^{15,16} The presence of constant visible coloration of the active media was not observed for a 5-h irradiation of the crystal by the VUV F_2 laser and with a 2-Hz repetition rate. Despite the presence of saturation processes in the active media, the maximum output energy at 172 nm from the excited crystal was 0.4 mJ for a nonoptimized optical cavity. The low-threshold energy flux and the linear part of Fig. 5 indicate the presence of high gain in the crystal.

CONCLUSIONS

Vacuum-ultraviolet (VUV) laser radiation has been obtained from a solid-state $\text{LaF}_3:\text{Nd}^{3+}$ laser at 172 nm

pumped by a molecular F_2 laser at 157 nm. The system exhibits an extremely high gain. Therefore good prospects become apparent for further developments of solid-state VUV tunable lasers based on these crystals with F_2 laser pumping. Studies of other Nd^{3+} -activated, as well as other rare-earth-ion-activated fluoride single crystals with the same pumping for UV and VUV laser development are in progress.

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